## Power Calculations

## Lecture 6 review:

- Power Calculations
- Starting your car (model)


## Today: (5.1, 5.2)

- Ideal inductors and capacitors
- Energy storage
- Practical capacitors and inductors

Why?

- To find out how much power is being delivered to (from) some device, e.g. loud speaker.
- Alternative, it may be undesirable to delivered power to some part of the circuit, because of the heat it generates.
- Basic idea is addressed in the first lecture.

How?
Use the Associated Reference Direction convention.


## Power Definitions

- $\mathrm{P}=\mathrm{VI}>0$ corresponds to the element absorbing power
- How can a circuit element absorb power?
- By converting electrical energy into heat (resistors in toasters), light (light bulbs), acoustic energy (speakers); by storing energy (charging a battery)
- Negative power - releasing power to the rest of the circuit.


## Conservation of Power

- Sum of the power absorbed by all circuit element must be zero.
- Concept: circuit elements are used to model all modes of energy conversion (heat, sound, batteries, voltage generators, etc.)
- Simple example:

$$
I=-2 \mathrm{~mA}
$$

Power released ( $\mathrm{VI}<0$ ) by the element on the left equals to the power absorbed by the element on the right.


## Example of Power Flowing into Current Source

What is the power flow into the current source in the circuit on the right?

Put an imaginary box enclosing the current
 source and apply the associated reference direction (ARD) to the current source.
(1) Assume $V$ convention on the right, ARD dictates that current has to go into + side of the terminals.
By KVL: $V=100 \mathrm{~mA} \times(60 \Omega+40 \Omega)=10 \mathrm{~V}$ By KCL: I =-100mA
The power entering the box (in this case, current source) is: Power $=\mathrm{VI}=10 \mathrm{~V} \times-100 \mathrm{~mA}=-1 \mathrm{~W} \quad \therefore 1 \mathrm{~W}$ is leaving the box


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## Example of Power Flowing into resistor

Find the power entering the $40 \Omega$ resistor as shown in the circuit.

By associate reference direction convention. the V and I are defined.

By inspection $I=100 \mathrm{~mA}$

$$
V=I \times 40 \Omega=4 V
$$

$\mathrm{P}=\mathrm{I} \times \mathrm{V}=100 \mathrm{~mA} \times 4 \mathrm{~V}=0.4 \mathrm{~W}$

$\therefore 0.4 \mathrm{~W}$ is entering the box (or $40 \Omega$ resistor).
(Using the similar strategy, it is found that 0.6 W is entering the $60 \Omega$ resistor.)
${ }^{\text {mesursmis } 200}$ Experimental Measurement of Power using ${ }^{\text {citai }}$ a Volt Source and a Ammeter
The box in the right contains an unknown circuit, experimentally it is found that the I-V diagram as shown.
(how? Set the voltage, and measure the current by the ammeter)


| V (set) | I (measured) | $\mathrm{P}=\mathrm{V} \mathrm{I}$ (compute) |
| :--- | :--- | :--- |
| -3 V | -18 mA | +54 mW (power entering) |
| -2 V | -16 mA | +32 mW |
| -1 V | -14 mA | +14 mW |
| 0 | -10 mA | 0 (no power transfer) |
| 1 V | -2 mA | -2 mW (power leaving) |
| 2 V | +20 mA | +40 mW |
| 3 V | +400 mA | +1.2 W |

$\therefore$ In the first and the third quadrant of the I-V curve, power is entering. In the second and the fourth quadrant, power is leaving.

## Instantaneous Power and Average Power

When $V$ and $I$ are function of time, we write as $v(t)$ and $i(t)$, then instantaneous power entering the box is:

$$
P(t)=v(t) i(t) \text { (instantaneous power, power as function of } \mathrm{t} \text { ) }
$$

This is particularly useful to compute the maximum power received or delivered at any instant.

$$
P_{A V}=\frac{1}{T} \int_{0}^{T} v(t) i(t) d t \quad \text { (time-averaged power) }
$$

This is particularly useful for computing the trend, the average power received or delivered.

Does DC current have instantaneous or time-average power?

## Special cases

If a voltage exists across a resistor R ,
Power dissipated in the resistor $=\mathrm{VI}=\mathrm{V}(\mathrm{V} / \mathrm{R})=\mathrm{V}^{2} / \mathrm{R}$ (Ohm's law applies in this case)

If a current flow, I, through a resistor $R$,
Power dissipated in the resistor $=\mathrm{VI}=(\mathrm{IR}) \mathrm{I}=I^{2} \mathrm{R}$ (Ohm's law also applies in this case)

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## Example (Time-averaged Power Transferred)




Find the timétaveraged power transferked into resistor R, if the time varying voltage source $v(t)$ is as shown on the right.

Observation: (1) $\mathrm{v}(\mathrm{t})$ is a square wave ( max at $\mathrm{V}_{0}$, min at 0 )
$i(t)$ is a $\qquad$ wave (max at , min at )
Instantaneous power $=v(t) i(t)=\quad$ or
(2) Period from 0 ms to 3 ms , or 1 ms to 4 ms , or 2 ms to 5 ms .
$P_{A V}=\frac{1}{0.003-0} \int_{0}^{0.003} v(t) i(t) d t$
$P_{A V}=\frac{1}{0.003}\left(\int_{0}^{0.001} 0 d t+\int_{0.001}^{0.002} \frac{V_{0}^{2}}{R} d t+\int_{0.002}^{0.003} 0 d t\right)=\frac{1}{0.003} \frac{V_{0}^{2}}{R}(0.001)=\frac{V_{0}^{2}}{3 R}$

## Multi-terminal Elements

We dealt with circuit element with 2 terminals in the preceding sections.

Here we have 4 terminals
$P=V_{1} I_{1}+V_{2} I_{2}+V_{3} I_{3}+V_{4} I_{4}$
Notice the reference direction
 of $I$ is toward the terminals, the

Show that the 2 terminals circuit and formula is a special case to our 4 terminals circuit formula.

Is $\mathrm{P}=\mathrm{VI}$ a special case of
$\mathrm{P}=\mathrm{V}_{1} \mathrm{I}_{1}+\mathrm{V}_{2} \mathrm{I}_{2}+\mathrm{V}_{3} \mathrm{I}_{3}+\mathrm{V}_{4} \mathrm{I}_{4}$ ?

## Example

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Generalize into a 2 terminals circuit:
$P=V_{1} I_{1}+V_{2} I_{2}$
Comparing the 2 circuits,
$I=l_{1}, \quad I=-1_{2}$
$\mathrm{P}=\mathrm{V}_{1} \mathrm{I}_{1}+\mathrm{V}_{2} \mathrm{I}_{2}$ becomes $\mathrm{P}=\mathrm{V}_{1} \mathrm{I}-\mathrm{V}_{2} \mathrm{I}=\left(\mathrm{V}_{1}-\mathrm{V}_{2}\right) \mathrm{I}$


But $\mathrm{V}=\mathrm{V}_{1}-\mathrm{V}_{2}$ (comparing the 2 circuits)
$\mathrm{P}=\mathrm{VI}$ which is our original Power formula.
$\therefore \mathrm{P}=\mathrm{VI}$ is a special case of $\mathrm{P}=\mathrm{V}_{1} \mathrm{I}_{1}+\mathrm{V}_{2} \mathrm{I}_{2}+\mathrm{V}_{3} \mathrm{I}_{3}+\mathrm{V}_{4} \mathrm{I}_{4}$ general formula.

## Battery model

We can model each battery as an ideal 12 V voltage source. But this would not help us to understand the difference between eight 1.5 V batteries in series and a 12 V car battery.

Next, we can model the battery (2 terminal voltage source) as a Thévenin equivalent circuit as shown below.

Both batteries have an identical $\mathrm{V}_{\mathrm{T}}$ (open circuit voltage), so what is the difference between the two?

By voltage divider formula: $V_{M}=V_{T} \frac{R_{M}}{R_{T}+R_{M}}$
Since $V_{T}$ and $R_{M}$ are the same for both batteries, the difference is in $\mathrm{R}_{\mathrm{T}}$.


The typical car battery has a Thévenin resistance of $0.05 \Omega$.

$$
V_{M}=12 \frac{0.4}{0.05+0.4}=10.67 \mathrm{~V} \quad I=\frac{12}{0.05+0.4}=26.67 \mathrm{~A}
$$

The typical AA size battery has a Thévenin resistance of $20 \Omega$.

$$
V_{M}=12 \frac{0.4}{20+0.4}=0.235 \mathrm{~V} . \quad I=\frac{12}{20+0.4}=0.588 \mathrm{~A}
$$

That is not sufficient voltage and current to start a car.
This is no surprising because you don't expect AA size batteries to be sufficient to start a car!
Typically, the resistance $\left(R_{T}\right)$ is a function of the cross section of the device (in this case, battery). The larger the cross section, the smaller the resistance.

## Duality

-     +         - 
- black, white
- Voltage, Current
- Capacitors, inductors.
- Power flow can be calculated from the expressions $\mathbf{P}=\mathbf{V I}$ for two-terminal circuit elements and $\mathbf{P}=\boldsymbol{\Sigma} \boldsymbol{V}_{n} \boldsymbol{I}_{\boldsymbol{n}}$ for multi-terminal circuit element. However, it is essential that the signs of the various voltages and currents be stated correctly.
- If voltage and current vary, the quantity $v(t) i(t)$ is known as the instantaneous power. The time-averaged power is the average over time of the instantaneous power.
- $\mathrm{Q}=\mathrm{CV}$
where $\mathrm{Q}=$ charge

$$
\begin{aligned}
& \text { V=voltage difference between } 2 \text { plates } \\
& \mathrm{C}=\text { capacitance }
\end{aligned}
$$

take derivative with respect to $t$ on both sides

- $d Q / d t=C d V / d t=i$
- $\mathrm{i}=\mathrm{CdV} / \mathrm{dt} \quad$ (remember I-V diagram)
- current = constant X time derivative of voltage
- Ohm's law tells us about the relationship between

V and I for a resistor. This equation describes the relationship between i and V for a capacitor.

## Capacitors (continue)

- The I-V relationship for a capacitor is:

$$
I_{A \rightarrow B}=C \frac{d}{d t}\left(V_{A B}\right)
$$



Where $C$ is the capacitance in Farad or $F, m F, \mu F, n F, p F$
Notice the current depends on the derivative. If the derivative is zero, then there is no current. The derivative is zero when the voltage remains constant and does not change with time.
An example would be: dc circuit.
$\therefore$ No current goes through a capacitor in a dc circuit.

## Capacitor example

- Find the $\mathrm{v}_{\mathrm{c}}(\mathrm{t})$ across the capacitor as shown. The current $I_{0}$ through the current source is constant.
Apply the I-V equation for capacitor from the previous page (when the current direction $A->B$, then Voltage is $\mathrm{V}_{\mathrm{AB}}$ )
$I_{A \rightarrow B}=C \frac{d}{d t}\left(V_{A B}\right)$


In the circuit on the right, current $\mathrm{I}_{0}$ is entering cap.
$I_{0}=C \frac{d}{d t}\left(V_{c}\right) \Rightarrow \frac{I_{0}}{C} d t=d V_{c} \Rightarrow V_{c}=\int d V_{c}=\frac{I_{0}}{C} \int d t=\frac{I_{0} t}{C}+K$
$\therefore$ The voltage is increasing proportional to time, cap. Is charged.

## Capacitor example

- Find the current $\mathrm{I}_{1}(\mathrm{t})$ that passes through the capacitor as shown. The voltage source is a sinusoid $V_{0} \sin \omega t$, where $\mathrm{V}_{0}$ and $\omega$ are given constants and t is time.

Since the voltage source is sinusoidal (change with time), the current across the capacitor is nonzero.

From circuit

$$
V_{A}-V_{B}=V_{0} \sin \omega t
$$



From previous page $\quad I_{A \rightarrow B}=C \frac{d}{d t}\left(V_{A B}\right)$
$I_{A \rightarrow B}=C \frac{d}{d t}\left(V_{0} \sin \omega t\right)$
$I_{A \rightarrow B}=C V_{0} \omega \cos \omega t$
But $i$ is from $B$ to $A$ direction:

$$
I_{B \rightarrow A}=I_{1}=-C V_{0} \omega \cos \omega t
$$

## ${ }^{\text {Hosaspmemen nemen }}$ Capacitors in series and parallel

- We studied that when $R_{1}$ and $R_{2}$ are in series, the $R_{\text {eq }}$ is equal to $R_{1}+R_{2}$, the $R_{\text {eq }}$ for resistor in parallel is:
$R_{1} R_{2} /\left(R_{1}+R_{2}\right)$
- What about capacitors?

If 2 capacitors are in parallel, the voltage across both would be the same $(=\mathrm{V})$

$$
I=C_{1} \frac{d V}{d t}+C_{2} \frac{d V}{d t}=\left(C_{1}+C_{2}\right) \frac{d V}{d t}=C_{e q} \frac{d V}{d t} \quad C_{e q}=C_{1}+C_{2}
$$

If 2 capacitors are in series, the current across both would be the same (=I)

$$
I=C_{1} \frac{d V_{1}}{d t}+C_{2} \frac{d V_{2}}{d t}
$$

## Capacitors in series

If 2 capacitors are in series, the current across both would be the same (=I) $\quad I=C_{1} \frac{d V_{1}}{d t}=C_{2} \frac{d V_{2}}{d t}$
In series, $\quad V=V_{1}+V_{2} \Rightarrow \frac{d V}{d t}=\frac{d V_{1}}{d t}+\frac{d V_{2}}{d t}$
From eq (1) $I=C_{1} \frac{d V_{1}}{d t}=C_{2} \frac{d V_{2}}{d t} \Rightarrow \frac{d V_{2}}{d t}=\frac{C_{1}}{C_{2}} \frac{d V_{1}}{d t} \quad$ Substitute into eq (2)

$$
\begin{equation*}
\Rightarrow \frac{d V}{d t}=\frac{d V_{1}}{d t}+\frac{C_{1}}{C_{2}} \frac{d V_{1}}{d t}=\frac{d V_{1}}{d t}\left(\frac{C_{2}+C_{1}}{C_{2}}\right) \tag{3}
\end{equation*}
$$

From eq (1) $I=C_{1} \frac{d V_{1}}{d t} \Rightarrow \frac{d V_{1}}{d t}=\frac{I}{C_{1}} \quad$ Substitute into eq (3)
$\Rightarrow \frac{d V}{d t}=\frac{I}{C_{1}}\left(\frac{C_{2}+C_{1}}{C_{2}}\right)=I\left(\frac{C_{2}+C_{1}}{C_{1} C_{2}}\right) \Rightarrow I=C_{\text {series }} \frac{d V}{d t} \quad \therefore C_{\text {series }}=\frac{C_{1} C_{2}}{C_{1}+C_{2}}$
Series capacitors are similar to parallel resistors.

## Inductors (example)

- Assume there is no current going thru an inductor, at time $=0$, a time varying current $\mathrm{i}(\mathrm{t})$ is applied thru the inductor terminals. What is the voltage across the inductor terminals as a function of time.

Apply the inductor equation in the previous page:

$$
v(t)=L \frac{d}{d t} i(t)
$$

If a time varying voltage $v(t)$ is applied across its terminals. What is the current thru the inductor as a function of time. Again, apply the inductor equation in the previous page:
$v(t)=L \frac{d}{d t} i(t) \Rightarrow d i(t)=\frac{1}{L} v(t) d t \Rightarrow i(t)=\int d i(t)=\int_{0}^{t} \frac{1}{L} v(t) d t$

## Inductors

- Ideal inductor is a 2-terminal device.

$$
\begin{equation*}
V_{A B}=L \frac{d}{d t} I_{A \rightarrow B} \tag{1}
\end{equation*}
$$

- Where L is a constant called inductance with unit in Henry or $\mathrm{H}, \mathrm{mH}, \mu \mathrm{H}, \mathrm{nH}$.


Notice similarity with capacitance equation $\quad I_{A \rightarrow B}$
$I_{A \rightarrow B}=C \frac{d}{d t}\left(V_{A B}\right)$

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- Write a loop equations for the loop current $i(t)$

The voltage drop in the inductor is: $L \frac{d}{d t} i(t)$
The voltage drop in the resistor is: $i(t) R$


So, the loop equation (KVL) is: $\quad L \frac{d}{d t} i(t)+i(t) R=0$
What if you have the same circuit in series with a capacitor C ?
Recall the I-V eq. for Capacitor: $i(t)=C \frac{d}{d t} v(t)$
$i(t)=C \frac{d}{d t} v(t) \Rightarrow d v(t)=\frac{1}{C} i(t) d t \Rightarrow v(t)=\int \frac{1}{C} i(t) d t$
The voltage drop in the capacitor is: $v(t)=\int \frac{1}{C} i(t) d t$
So, the new loop equation (KVL) is: $\quad L \frac{d}{d t} i(t)+i(t) R+\int \frac{1}{C} i(t) d t=0{ }_{28}$

## Parallel/Series Inductors

- Inductor in series is similar to resistor (sum):

$$
L_{\text {series }}=L_{1}+L_{2}
$$

- Inductor in parallel is similar to resistor (product over sum):

$$
L_{\text {parallel }}=\frac{L_{1} L_{2}}{L_{1}+L_{2}}
$$

## Energy Storage (continue)

$$
E=\int_{0}^{t} p(t) d t=\int_{0}^{t} v(t) i(t) d t
$$

Recall the I-V eq. for Capacitor: $\quad i(t)=C \frac{d}{d t} v(t)$
$E=\int_{0}^{t} v i d t=\int_{0}^{t} v C \frac{d v}{d t} d t=\int_{0}^{t} \frac{1}{2} C d\left(v^{2}\right)$
$E=\frac{1}{2} C[v(t)]^{2}-\frac{1}{2} C[v(0)]^{2}$
If the Capacitor is initially uncharged:
$E=\frac{1}{2} C[v(t)]^{2}$

$\therefore$ Energy store in a capacitor is: $1 / 2 \mathrm{CV}^{2}$

## Energy Storage

|  | Element | equation | energy+ or - ? | energy eq. |
| :--- | :--- | :--- | :--- | :--- |
| - Resistor R | $\mathrm{V}=\mathrm{IR}$ | dissipate energy | $\mathrm{V}^{2} / \mathrm{R}$ or $\mathrm{I}^{2} \mathrm{R}$ |  |
| - Capacitor C | $\mathrm{I}=\mathrm{CdV} / \mathrm{dt}$ | stored | ??? |  |
| - Inductor L | $\mathrm{V}=\mathrm{LdI} / \mathrm{dt}$ | stored | ??? |  |

Assume the capacitor is uncharged, at $t=0$, a voltage $v(t)$ is applied. The instantaneous power enter the capacitor is: $p(t)=v(t) i(t)$

The energy enter the capacitor (from time $=0$ to t)is:

$$
E=\int_{0}^{t} p(t) d t=\int_{0}^{t} v(t) i(t) d t
$$



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## Energy Storage (continue)

- Once energy is stored in capacitor, is there way we can regain the energy?

Consider the circuit on the right.
Suppose the capacitor is initially charged to voltage V , is to discharge to an external circuit.


The energy recovered from the capacitor (entered the external circuit) after an infinite length of time:

$$
\begin{aligned}
& E=\int_{0}^{\infty} p(t) d t=\int_{0}^{\infty} v_{1}(t) i_{1}(t) d t \\
& E=\int_{0}^{\infty} v(t)\left[-C \frac{d v}{d t}\right] d t=-C \int_{0}^{\infty} v(t) d v=-C \int_{0}^{\infty} d\left(\frac{v^{2}}{2}\right) \quad \begin{array}{l}
\text { Minus sign is due to the } \\
\text { reference direction }
\end{array} \\
& E=-\frac{1}{2} C\left[v(\infty)^{2}-v(0)^{2}\right] \quad \text { The voltage is fully discharged }(\mathrm{V}=0) \text { when } \mathrm{t}=\infty \\
& E=\frac{1}{2} C\left[v(0)^{2}\right]
\end{aligned}
$$

## Energy storage for inductors

$$
\begin{aligned}
& E=\int_{0}^{t} v(t) i(t) d t=\int_{0}^{t} L \frac{d i}{d t} i d t=\int_{0}^{t} \frac{1}{2} L d\left(i^{2}\right) \\
& E=\frac{1}{2} L[i(t)]^{2}-\frac{1}{2} L[i(0)]^{2}
\end{aligned}
$$

$E=\frac{1}{2} L I^{2}$
Where I is the final current at time $t$, assumed the current through the inductor is zero.

## Practical Capacitors and inductors

Practical capacitor = ideal capacitor in series with a resistor


The resistor part dissipates energy, thus, practical capacitor can never retain energy definitely, e.g. every DRAM cell need to be refreshed periodically to retain its value.

Capacitors use below 1GHz: mica, ceramic, and tantalum (see Figure a). Capacitors are specified by their capacitance value, maximum voltage applied across terminals, their tolerance.

## Practical Capacitors and inductors (cont)

- a practical inductor can be replaced by an ideal inductor in series with a resistor and then in parallel with a capacitor

Again, practical inductor can dissipate energy because of the present of the resistor.


