

EE243 Advanced Electromagnetic Theory

Lec # 16 Dielectric Waveguides

- **Recap Solutions for Homework Set 6**
- **Review for the Midterm Examination**
- **Dispersion Equation Dielectric Slab Waveguides**
- **Modes in Dielectric Slab Waveguides**

Reading: Jackson 8.11 and Harrington 4.7.

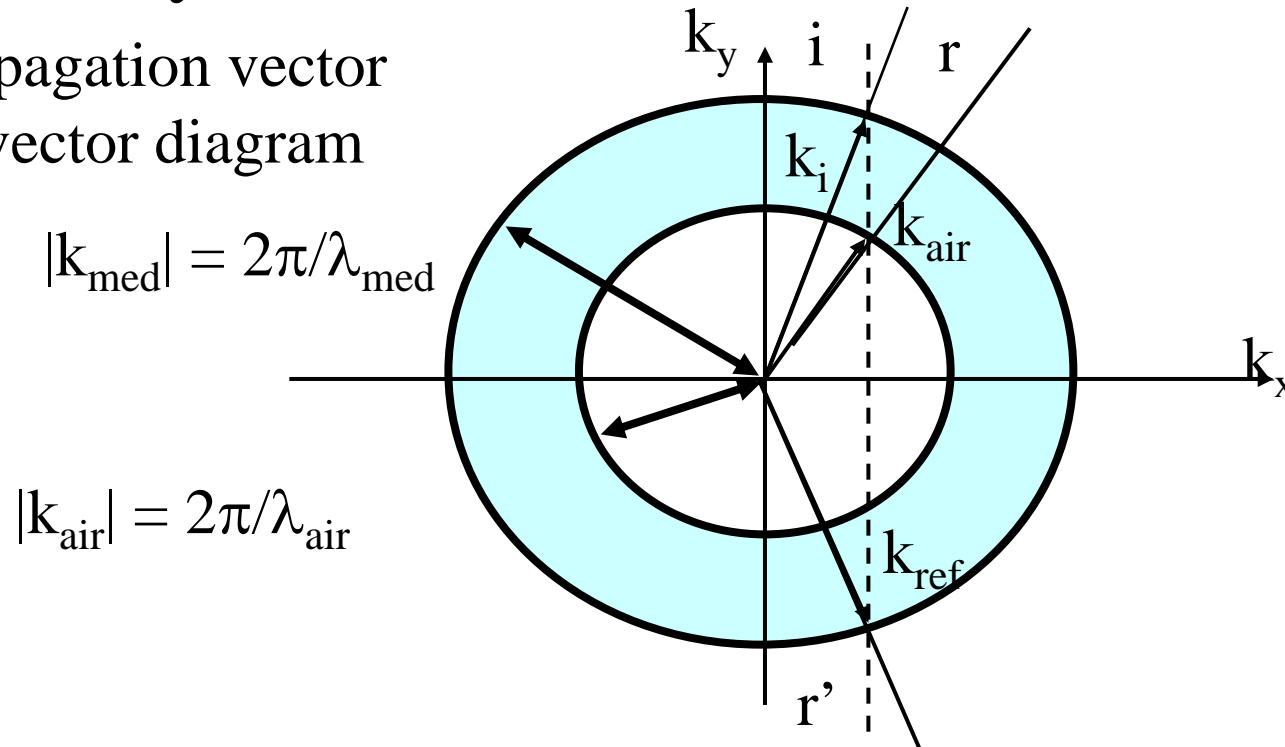
Midterm Exam

Review Th, Oct 19

- In Class Tuesday October 24th
- Covers material through Chapter 7 (Lecture 12)
- Open Book, Open Notes, Bring Calculator, Paper Provided
- Topics
 - Green's functions free space and use in Theorems and concepts with emphasis on statics
 - Separation of variables in rectangular coordinates using N-1 and N method
 - Time-Harmonic ME, planewaves, boundary conditions, and dispersion

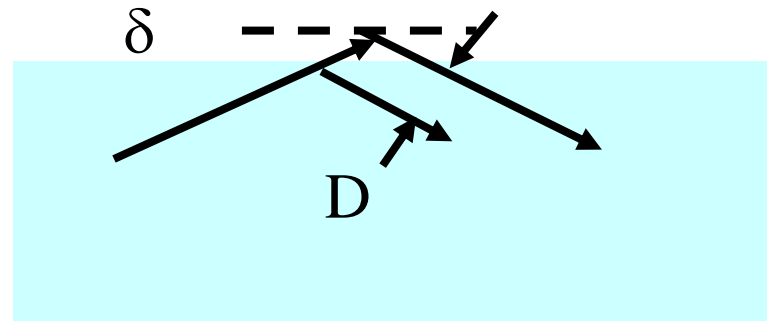
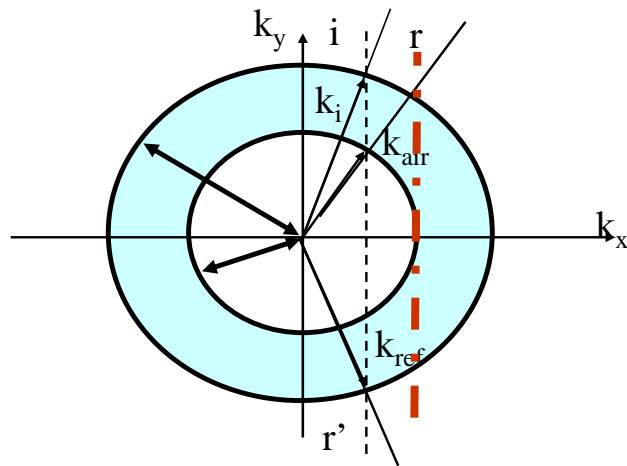
Physical Effects: Wave Direction Change

Propagation vector
k-vector diagram



- Draw concentric circles of radius k_{air} and k_{med}
- Incident wave has k vector given (arrow k_i)
- Find the component parallel to the surface (dotted line)
- Force the k-vector in air k_{air} and k-vector reflected k_{ref} to have the same parallel component (lie on dotted line)
- Choose point on the circle to give these new k-vectors (arrows) the correct length for the wave equation in the media that they are in

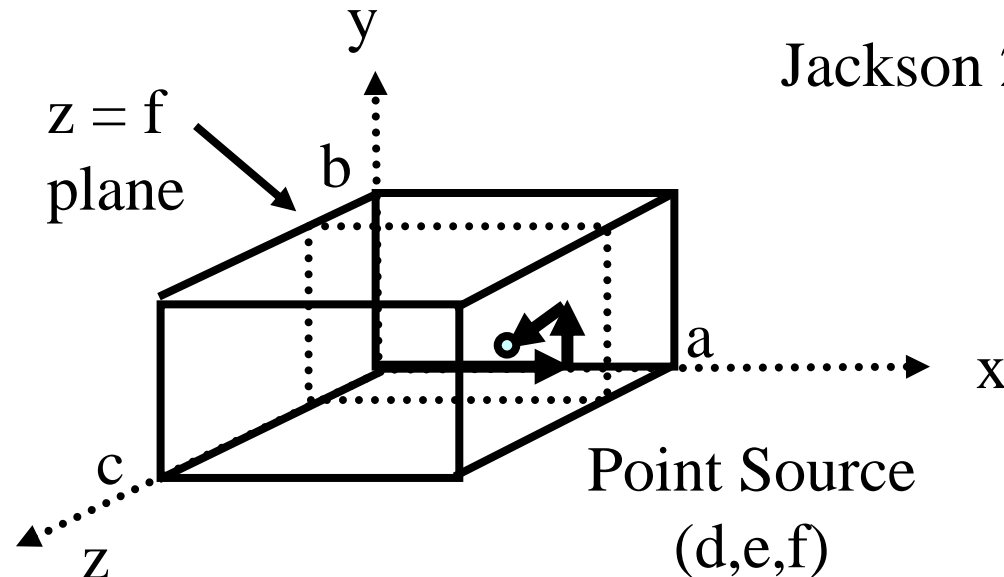
Plane Interface: Physical Effects



- Total Internal Reflection
 - Parallel part of $k_{\text{med}} > k_0$
- Brewster Angle
 - Polarization in plane of incidence reflection coefficient goes to zero giving 100% transmission
- Polarization dependent reflection phase change
 - Converts linear to part circular polarization
 - Beam energy penetration δ and spot shift D (Goos-Hanchen effect)

Separation of Variables: Source Strategy

Jackson 2.9

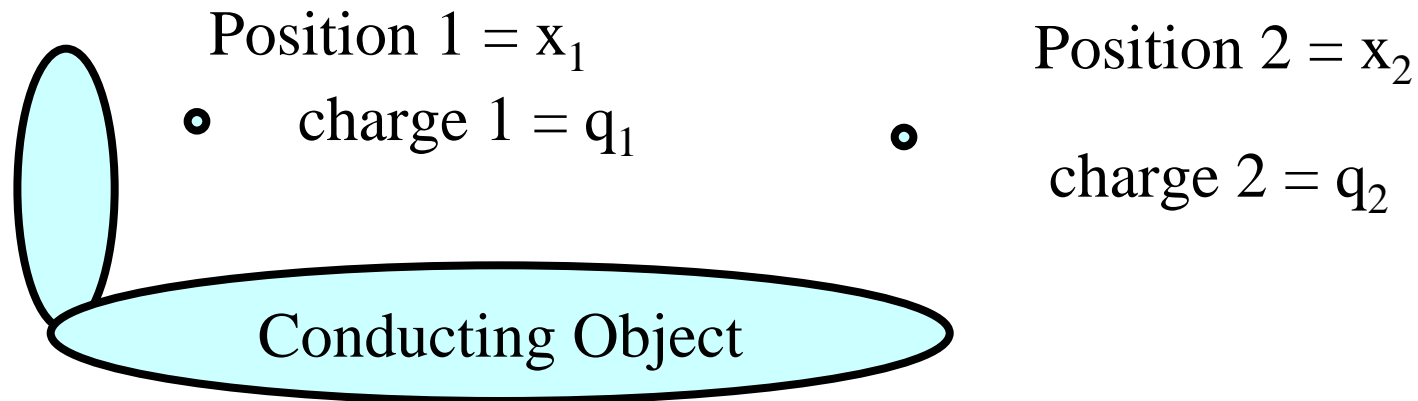


- View source as being on $z = f$ plane.
- Require $\Phi_2 - \Phi_1 = D(x, y) / \epsilon_0$ at $z = f$
- Also require at $z = f$ $(\bar{E}_2 - \bar{E}_1) \cdot \hat{n} = \sigma_{SURFACE}(x, y) / \epsilon_0$
- Multiply each of these equations by one of the composite eigenfunctions and integrate over x, y cross-section
- Gives two equations relating A_{nm} and B_{nm} for the same nm .

$$\nabla^2 \psi = -\frac{4\pi q}{\epsilon_0} \delta(x-d)\delta(y-e)\delta(z-f)$$

Reciprocity

Lecture 4



$$\Phi(x_2, q_1) = \Phi(x_1, q_2)$$

Proof:

- Green's Theorem
- Poisson's equation for $\Phi(x_2, q_1)$ and $\Phi(x_1, q_2)$ causes volume integral to give $\Phi(x_2, q_1) - \Phi(x_1, q_2)$
- In surface integral use homogeneous boundary condition to replace potential with derivative and integrand vanishes at every point on the boundary

Reciprocity for Green's Function in Jackson

$$G(\bar{x}, \bar{x}') = G(\bar{x}', \bar{x})$$

Integral Equation to Find Surface Charge

$$\Phi(\bar{x}) = \frac{1}{4\pi\epsilon_0} \int_V \rho(\bar{x}') G(\bar{x}, \bar{x}') d^3 x' + \frac{1}{4\pi} \oint_S \left[G(\bar{x}, \bar{x}') \frac{\delta\Phi}{\delta n'} - \Phi(\bar{x}') \frac{\delta G(\bar{x}, \bar{x}')}{\delta n'} \right] da'$$

- Example: Grounded Conducting Object and $\rho(x)$

$\Phi = 0 \Rightarrow$ all of the F surface term drop out

Lecture 5

$d\Phi/dn' = \sigma_{\text{surface}}$ remains

Since Φ is known at every point on object restrict x to be on the object

Gives and integral equation for the surface charge

$$0 = \left[\frac{1}{4\pi\epsilon_0} \int_V \rho(\bar{x}') G(\bar{x}, \bar{x}') d^3 x' + \frac{1}{4\pi} \oint_S \left[G(\bar{x}, \bar{x}') \frac{\delta\Phi}{\delta n'} \right] da' \right]_{\bar{x}_{\text{on_object}}}$$

Generally the Green's function for free space is used

$$0 = \left[\frac{1}{4\pi\epsilon_0} \int_V \rho(\bar{x}') \frac{1}{|\bar{x} - \bar{x}'|} d^3 x' + \frac{1}{4\pi} \oint_S \left[\frac{1}{|\bar{x} - \bar{x}'|} \frac{\delta\Phi}{\delta n'} \right] da' \right]_{\bar{x}_{\text{on_object}}}$$

Overview

- Source free guided wave solutions can exist on dielectric slabs, layers and fibers.
- The necessary conditions for their longitudinal propagation constant are found by representing the fields and matching boundary conditions on their transverse field to determine eigenvalues.
- The transverse behavior is exponential outside the dielectric and oscillatory inside the dielectric.
- The physical characteristics on the modes are quite similar to those in metal waveguides and include TE/TM classification, orthogonality, cut-off, etc.

Dielectric Waveguides

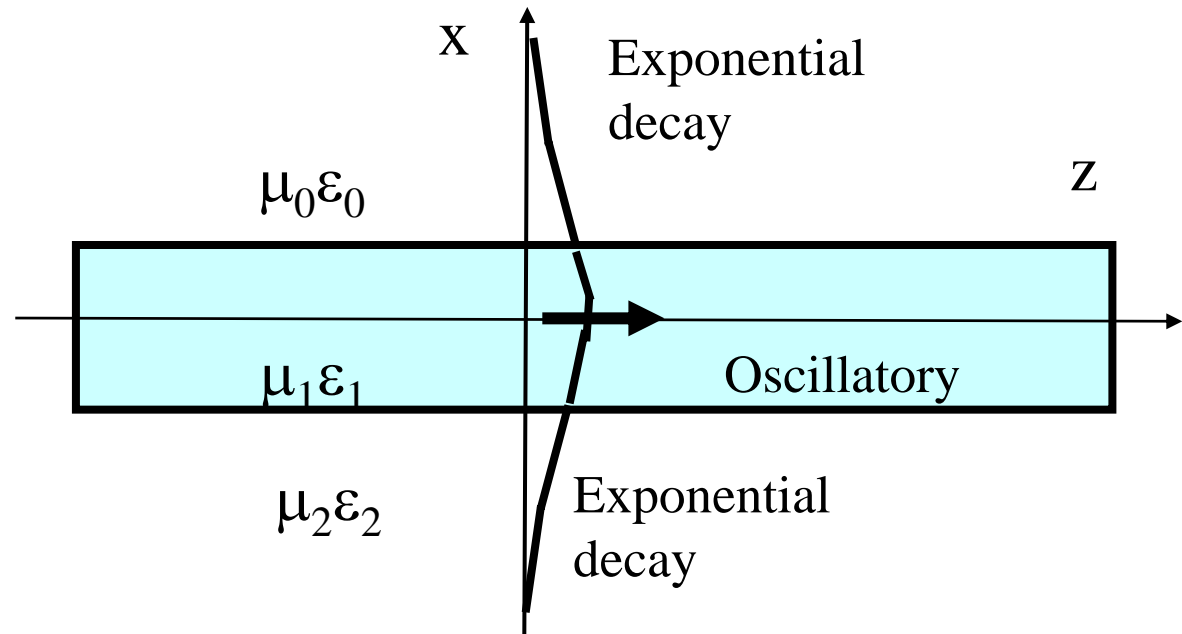
$$e^{j\omega t}$$

$$e^{-jk_z z}$$

$$v_0 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_0}$$

$$k_x = \sqrt{\omega^2 \mu_0 \epsilon_1 - k_z^2}$$

$$v_2 = \sqrt{k_z^2 - \omega^2 \mu_2 \epsilon_2}$$



- Three regions
- Choose TM (or TE)
- Will have H_y , E_z and E_x (E_y , H_x , and H_z)

Dielectric Waveguides

$$H_{0y}^+(\bar{x}) = H_0^+ \hat{y} e^{-\nu_0 x} e^{-jk_z z}$$

$$H_{1y}(\bar{x}) = H_1^+ \hat{y} e^{-jk_x x} e^{-jk_z z} + H_1^- \hat{y} e^{+jk_x x} e^{-jk_z z}$$

$$H_{2y}^-(\bar{x}) = H_2^- \hat{y} e^{+\nu_0 x} e^{-jk_z z}$$

- Consider TM w/r z case
- Write expression for H_y in each of three regions (above, in and below dielectric).
- Note: Include Kinetic boundary condition in expressions

Dielectric Waveguides

$$\nabla \times \bar{H} = j\omega\epsilon\bar{E}$$

$$E_z(\bar{x}) = \frac{-1}{j\omega\epsilon} \frac{\partial H_y(\bar{x})}{\partial x}$$

$$E_{0z}^+(\bar{x}) = \frac{+v_0}{j\omega\epsilon} H_{0y}^+(\bar{x})$$

$$E_{1z}^+(\bar{x}) = \frac{-jk_x}{j\omega\epsilon} H_{1y}^+(\bar{x}) + \frac{jk_x}{j\omega\epsilon} H_{1y}^-(\bar{x})$$

$$E_{2z}^-(\bar{x}) = \frac{-v_2}{j\omega\epsilon} H_{2y}^+(\bar{x})$$

- Find E_z in each of three regions
- Apply dynamic boundary conditions (four)
 - H_y continuous at top and bottom of dielectric
 - E_z continuous at top and bottom of dielectric

Dielectric Waveguide: Dispersion Eq.

Harrington 4.7 Special case of air on top and bottom, thickness a

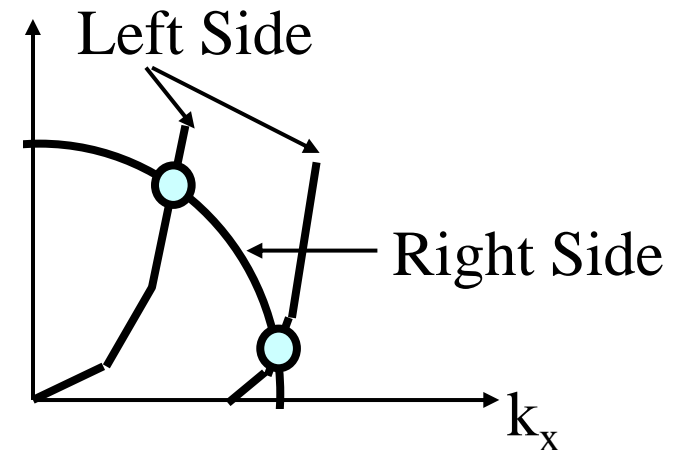
	TM	TE
odd	$\frac{k_x a}{2} \tan \frac{k_{x0} a}{2} = \frac{\epsilon_1}{\epsilon_0} \frac{v_0 a}{2}$	$\frac{k_x a}{2} \tan \frac{k_{x0} a}{2} = \frac{\mu_1}{\mu_0} \frac{v_0 a}{2}$
even	$-\frac{k_x a}{2} \cot \frac{k_{x0} a}{2} = \frac{\epsilon_1}{\epsilon_0} \frac{v_0 a}{2}$	$-\frac{k_x a}{2} \cot \frac{k_{x0} a}{2} = \frac{\mu_1}{\mu_0} \frac{v_0 a}{2}$

- Convenient within TM and TE to distinguish between even ($\cos k_x x$) and odd ($\sin(k_y x)$) variations
- Results in four dispersion relationships
 - Two for TM
 - Two for TE

Dielectric Waveguide: Physical Nature

Harrington 4.7 Special case of air on top and bottom, thickness a

$$\text{odd TM} \quad \frac{k_x a}{2} \tan \frac{k_{x0} a}{2} = \frac{\epsilon_1}{\epsilon_0} \frac{v_0 a}{2}$$



- Right hand side is a circle; Left hand side is spikes in \tan (See H Fig 4-11)
- Odd $\sin(k_y x)$ variations have no cut-off (always exist) in both TM and TE
- Multiple solutions (intersections) give multiple modes
- Additional new mode about every half wavelength of oscillatory variation.
- Weighted by material contrast $\sqrt{\mu_1 \epsilon_1 - \mu_0 \epsilon_0}$

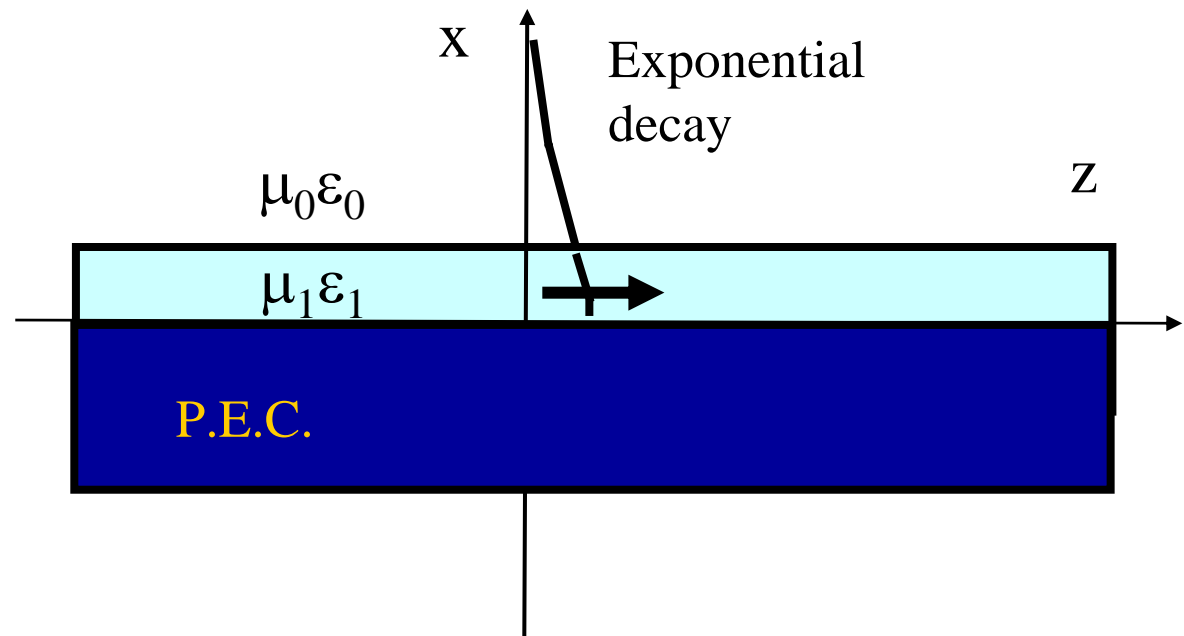
Surface-Guided Waves

$$e^{j\omega t}$$

$$e^{-jk_z z}$$

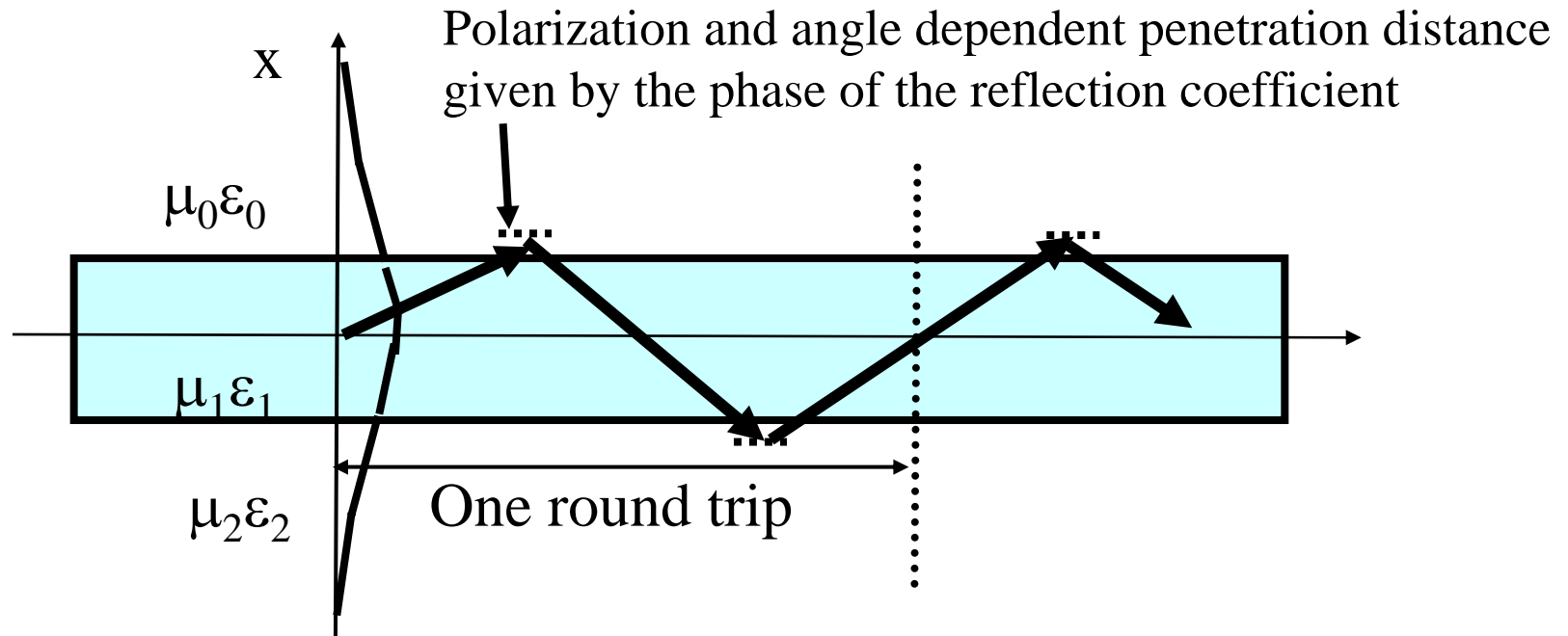
$$v_0 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_0}$$

$$k_x = \sqrt{\omega^2 \mu_0 \epsilon_1 - k_z^2}$$



- Two regions
- Choose TM (or TE)
- Will have half of the solutions from the symmetric dielectric slab: TM odd and TE even of the slab

Dielectric Waveguides: Resonance View



- Add up phase of transverse round trip = $n2\pi$
- Use the phase of the reflection coefficient to account for penetration of fields outside dielectric
- This phase will depend on polarization and angle and is thus must be found iteratively