

EE243 Advanced Electromagnetic Theory

Lec # 11: Plane Electromagnetic Waves

- **Plane Waves in a Nonconducting Medium**
- **Linear and Circular Polarization**
- **k-space view of waves in media**
- **Reflection and Refraction at Plane Interfaces**
- **Physical Phenomena Associated with Reflection**

Reading: Jackson Ch 7.1-7.5 (skip 7.6 and 7.7)

Overview

- In a source free region for time-harmonic ($e^{-j\omega t}$) signals Maxwell's Equations can be reduced to two coupled curl equations.
- These two curl equations
 - Combine to produce the wave equation for E or H
 - The eigenfunctions for these wave equations are plane waves described by propagation direction vectors called k-vectors that have length $2\pi/\lambda$ and result in wave velocity c .
 - Have zero divergence and make the vectors E and H perpendicular to the direction of propagation.
 - Make E and H vectors perpendicular to each other
 - Are sufficient at material boundaries to
 - require the components of the k-vector parallel to the surface to be the same on both sides of the boundary (Kinematic B.C.)
 - Require tangential E and H continuous; normal D and B continuous at the boundary (Dynamic B.C.)

Time-Harmonic Maxwell Equations

Time-Varying

$$\nabla \cdot \bar{D} = \rho$$

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}$$

$$\nabla \cdot \bar{B} = 0$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$$

Assume
No Sources

$$e^{-i\omega t}$$

$$\bar{J} = 0$$

$$\rho = 0$$

$$\bar{D} = \epsilon \bar{E}$$

$$\bar{B} = \mu \bar{H}$$

Time-Harmonic
Source Free

$$\nabla \times \bar{E} - i\omega \bar{B} = 0$$

$$\nabla \times \bar{B} + i\omega \mu \epsilon \bar{E} = 0$$

Implicit

$$\nabla \cdot \bar{D} = \rho$$

$$\nabla \cdot \bar{B} = 0$$

Wave Equation: Derivation

$$\nabla \times \bar{E} - i\omega \bar{B} = 0$$

$$\nabla \times \nabla \times \bar{E} - i\omega \nabla \times \bar{B} = 0$$

$$\nabla(\nabla \cdot \bar{E}) - \nabla^2 \bar{E} - i\omega(-i\omega \mu \epsilon \bar{E}) = 0$$

$$0 - \nabla^2 \bar{E} - \omega^2 \mu \epsilon \bar{E} = 0$$

$$\nabla^2 \bar{E} + \omega^2 \mu \epsilon \bar{E} = 0$$

$$\nabla^2 \bar{B} + \omega^2 \mu \epsilon \bar{B} = 0$$

- Take curl of curl E Eq.

- Sub: for curl curl

- Sub for curl B

- Use Div E = 0

- Similar Eq for B

Wave Equation: Plane Wave Solution

$$\nabla^2 \bar{E} + \omega^2 \mu \epsilon \bar{E} = 0$$

$$\bar{E} = \bar{E}_0 e^{i\bar{k} \cdot \bar{x}}$$

$$\nabla \rightarrow i\bar{k}$$

$$\nabla \cdot \rightarrow i\bar{k} \cdot$$

$$\nabla \times \rightarrow i\bar{k} \times$$

$$\nabla^2 \rightarrow -\bar{k}^2$$

$$-\bar{k}^2 + \omega^2 \mu \epsilon \bar{E} = 0$$

$$|\bar{k}| = \omega \sqrt{\mu \epsilon} = \frac{\omega}{c} = \frac{2\pi}{\lambda}$$

- Use 3D Fourier Expansion type eigenfunction where the vector \mathbf{k} is the propagation vector called the \mathbf{k} -vector
- Differential operators become algebraic operators
- Wave equation gives a constraint on the length of the \mathbf{k} -vector
- The \mathbf{k} -vector is reciprocal to the space variation wavelength

Plane Wave: Vector Properties

$$\bar{E} = \bar{E}_0 e^{i\bar{k} \cdot \bar{x}}$$

$$|\bar{k}| = \omega \sqrt{\mu \epsilon} = \frac{\omega}{c} = \frac{2\pi}{\lambda}$$

$$\nabla \cdot \bar{E} = 0 \rightarrow i\bar{k} \cdot \bar{E} = 0$$

$$\nabla \cdot \bar{B} = 0 \rightarrow i\bar{k} \cdot \bar{B} = 0$$

$$\bar{B} = \frac{1}{i\omega} \nabla \times \bar{E} \rightarrow \bar{B} = \frac{1}{i\omega} i\bar{k} \times \bar{E}$$

$$\rightarrow \bar{B} = \sqrt{\mu \epsilon} \hat{k} \times \bar{E}$$

$$\rightarrow \bar{H} = \frac{1}{\sqrt{\mu / \epsilon}} \hat{k} \times \bar{E} = \frac{1}{Z_0} \hat{k} \times \bar{E}$$

$$Z_0 = \sqrt{\mu / \epsilon} = 377 \text{ Ohms}$$

- Start with a vector in 3D and variation 3D
- $\text{Div } E = 0 \Rightarrow k$ perpendicular to E
- $\text{Div } B = 0 \Rightarrow$ perpendicular to E
- Because E is perpendicular to k the fact that $B \sim k$ cross E then implies B is perpendicular to E
- That is all 3 (k, B, E) are perpendicular to each other and that there are no fields in the direction of propagation

Inhomogeneous Plane Waves $|\mathbf{k}|^2 > \omega^2\mu\varepsilon$

$$\bar{\mathbf{E}} = \bar{\mathbf{E}}_0 e^{i\bar{\mathbf{k}} \cdot \bar{\mathbf{x}}}$$

$$\bar{\mathbf{k}} = \bar{k}_r + i\bar{k}_i$$

$$k_0 = \omega\sqrt{\mu\varepsilon} = \frac{\omega}{c} = \frac{2\pi}{\lambda}$$

$$\bar{\mathbf{k}} \cdot \bar{\mathbf{k}} = k_0^2$$

$$\text{Re}(\bar{\mathbf{k}} \cdot \bar{\mathbf{k}}) = \bar{k}_r^2 - \bar{k}_i^2 = k_0^2$$

$$\text{Im}(\bar{\mathbf{k}} \cdot \bar{\mathbf{k}}) = 2\bar{k}_r \cdot \bar{k}_i = 0$$

Evanescent waves that stay near or surface and explain phenomena such as tunneling across gaps.

- The \mathbf{k} -vector can be a vector with complex components and the imaginary part can describe exponential attenuation
- The wave equation requires the dot product with itself to
 - have a real part $\omega^2\mu\varepsilon$
 - have the imaginary part perpendicular to the real
- Thus the direction of maximum attenuation must be perpendicular to the direction of propagation

Plane-Wave: Poynting's Theorem

$$S = \frac{1}{2} (\overline{E} \times \overline{H}^*) = \frac{1}{2} \sqrt{\frac{\mu}{\epsilon}} |E_0|^2$$

$$w_e = \frac{1}{4} (\overline{E} \cdot \overline{D}^*) = \frac{\epsilon}{4} (\overline{E} \cdot \overline{E}^*) = \frac{\epsilon}{4} |E_0|^2$$

$$w_m = \frac{1}{4} (\overline{B} \cdot \overline{H}^*) = \frac{1}{4\mu} (\overline{B} \cdot \overline{B}^*) = \frac{\mu\epsilon}{4\mu} |E_0|^2$$

$$u = \frac{1}{4} \left(\epsilon \overline{E} \cdot \overline{E}^* + \frac{1}{\mu} (\overline{B} \cdot \overline{B}^*) \right) = \frac{\epsilon}{2} |E_0|^2$$

- Plug in Vector E with complex numbers for components
- Poynting vector, stored energy is balanced

Linear and Circular Polarization

- Choose propagation in z direction and let E have components in x and y directions

$$\bar{E} = \bar{E}_0 e^{i\bar{k} \cdot \bar{x}}$$

$$\bar{E}_0 = E_{0x} \hat{x} + E_{0y} \hat{y} + E_{0z} \hat{z}$$

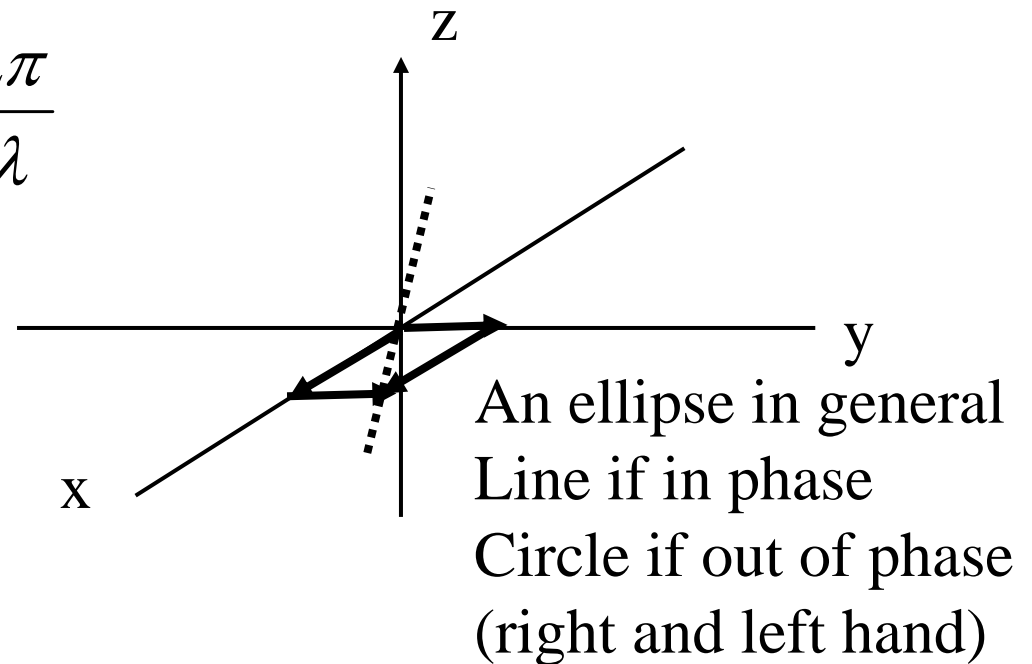
$$|\bar{k}| = k_0 = \omega \sqrt{\mu\epsilon} = \frac{\omega}{c} = \frac{2\pi}{\lambda}$$

$$\bar{k} = k_x \hat{x} + k_y \hat{y} + k_z \hat{z}$$

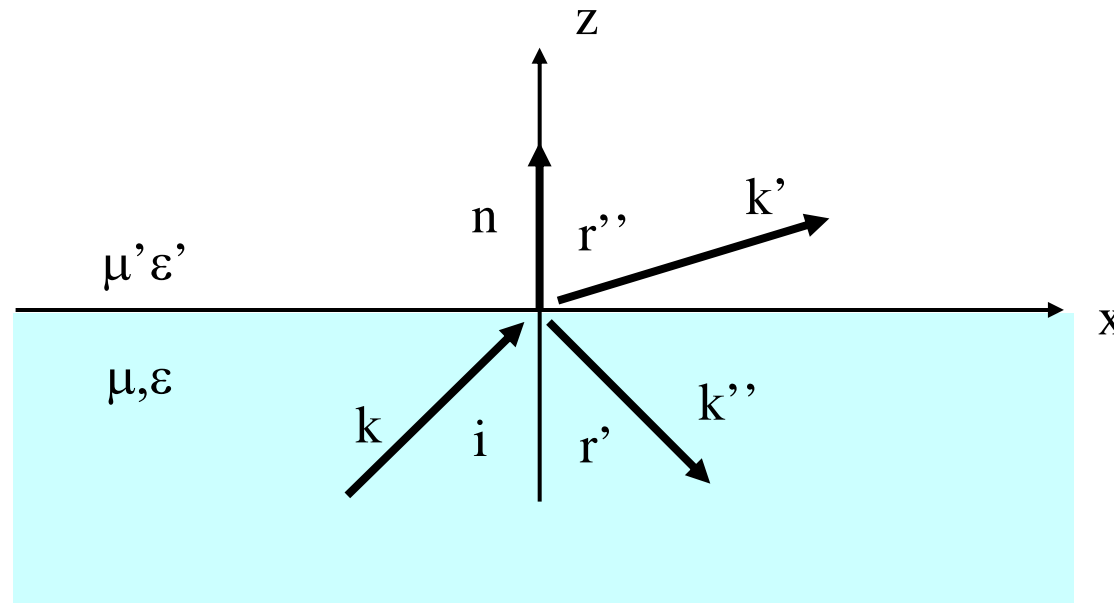
$$\bar{E}_0 = E_{0x} \hat{x} + E_{0y} \hat{y}$$

$$\bar{k} = k_0 \hat{z}$$

$$\bar{E}_0 = |E_{0x}| \cos(k_0 z - \omega t + \phi_x) \hat{x} + |E_{0y}| \cos(k_0 z - \omega t + \phi_y) \hat{y}$$



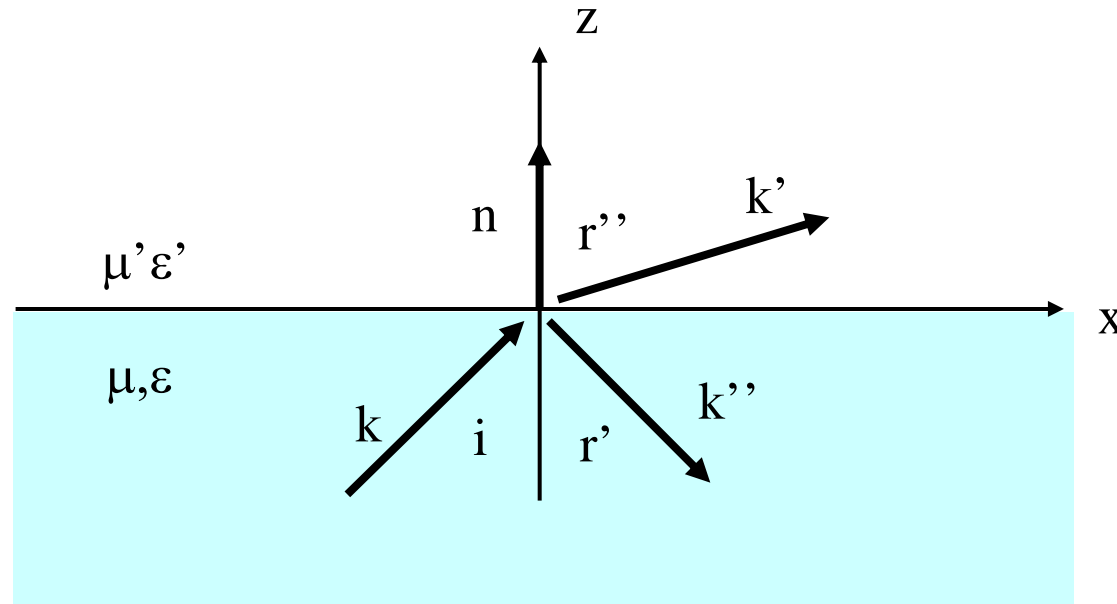
Reflection and Refraction at a Plane Interface



- Wave incident from below at angle i
- Generates transmitted (refracted) wave at angle r''
- Also generates reflected wave at angle r'

Note: The $z = 0$ plane where the boundary conditions are applied is for all x values and all y values

Plane Interface: Kinetic Boundary Conditions



- Since the $z = 0$ plane covers the range of the full Fourier representation each of the three waves must have the same eigenfunction variation along the boundary in x and y (i.e. the same k_x and k_y)

$$\left(\bar{\mathbf{k}} \cdot \bar{\mathbf{x}}\right)_{z=0} = \left(\bar{\mathbf{k}}' \cdot \bar{\mathbf{x}}\right)_{z=0} = \left(\bar{\mathbf{k}}'' \cdot \bar{\mathbf{x}}\right)_{z=0}$$

$$k_x = k'_x = k''_x \mapsto k_y = k'_y = k''_y$$

Plane Interface: Dynamic Boundary Conditions

- Wave Eq = 2nd order => 2 boundary conditions
- Two independent vector orientations => 4 boundary conditions total
- Choose 4 from 6 possible and express in terms of E
 - Normal D and B continuous (2)
 - Tangential E and H continuous (4)

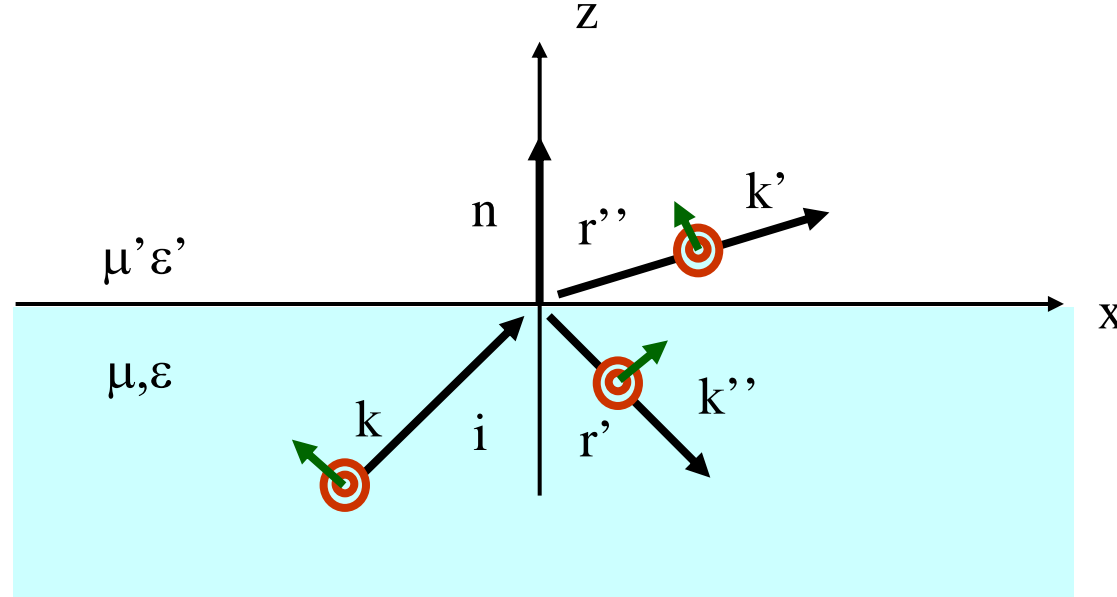
$$\left[\varepsilon \left((\bar{E}_0 + \bar{E}_0'') - \varepsilon' \bar{E}_0' \right) - \varepsilon' \bar{E}_0' \right] \cdot \hat{n} = 0$$

$$\left[\bar{k} \times \bar{E}_0 + \bar{k}'' \times \bar{E}_0'' - \bar{k}' \times \bar{E}_0' \right] \cdot \hat{n} = 0$$

$$\left(\bar{E}_0 + \bar{E}_0'' - \bar{E}_0' \right) \times \hat{n} = 0$$

$$\left[\frac{1}{\mu} \left(\bar{k} \times \bar{E}_0 + \bar{k}'' \times \bar{E}_0'' \right) - \frac{1}{\mu'} \bar{k}' \times \bar{E}_0' \right] \times \hat{n} = 0$$

Plane Interface: Solution



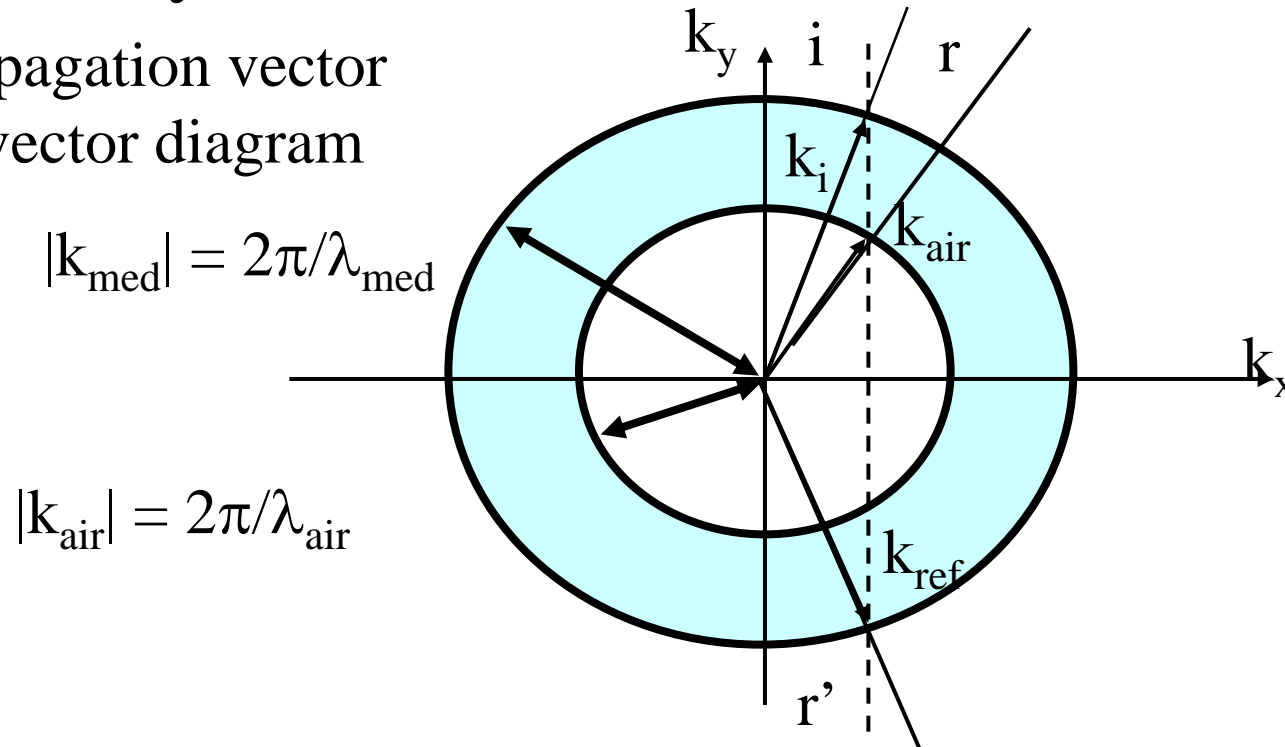
- Unknowns are the transmitted and reflected wave complex amplitudes **in the plane of incidence** and **perpendicular to the plane of incident**.
- These two vector orientations can be solved independently from each other
- See Jackson Page 305 for the detailed results

Plane Interface: Physical Effects

- Refraction
 - Wave direction change
- Total Internal Reflection
 - Only Evanescent fields outside
 - Tunneling across a gap
- Brewster Angle
 - 100% transmission
- Polarization dependent phase change
 - Converts linear to part circular polarization
 - Beam spot shift (Goos-Hanchen effect)

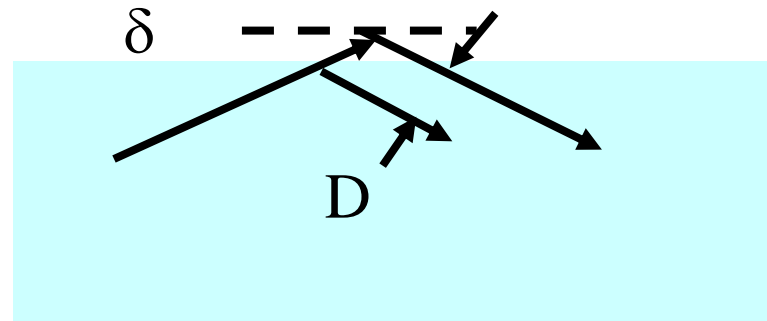
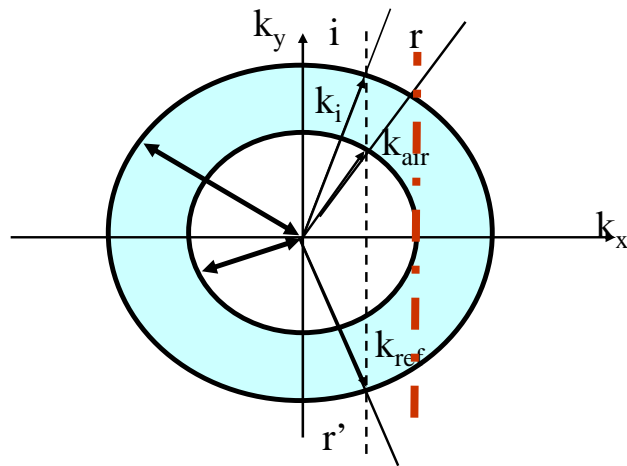
Physical Effects: Wave Direction Change

Propagation vector
k-vector diagram



- Draw concentric circles of radius k_{air} and k_{med}
- Incident wave has k vector given (arrow k_i)
- Find the component parallel to the surface (dotted line)
- Force the k-vector in air k_{air} and k-vector reflected k_{ref} to have the same parallel component (lie on dotted line)
- Choose point on the circle to give these new k-vectors (arrows) the correct length for the wave equation in the media that they are in

Plane Interface: Physical Effects



- Total Internal Reflection
 - Parallel part of $k_{med} > k_0$
- Brewster Angle
 - Polarization in plane of incidence reflection coefficient goes to zero giving 100% transmission
- Polarization dependent reflection phase change
 - Converts linear to part circular polarization
 - Beam energy penetration δ and spot shift D (Goos-Hanchen effect)