

EE243 Advanced Electromagnetic Theory

Lec # 9: Maxwell Equations

- **Maxwell Equations (including Faraday's EMF)**
- **Vector and Scalar Potentials**
- **Gauge Conditions to Give Potentials Useful Properties**
- **Green's Function for the Wave Equation**
- **Retarded Solutions**
- **Macroscopic Maxwell Equations Applications**

Reading: Jackson Ch 6.1-6.6 (lite on 6.5 and 6.6)

Overview

- Maxwell was able to integrate Faraday's observations into a single set of consistent equations for both statics and dynamics and Maxwell also considered light to be an electromagnetic phenomena.
- Key Ideas
 - Add displacement current
 - Re-derive the vector and scalar potential
 - Use Gauge conditions to tie down arbitrary nature
 - Use Fourier representation to find the time-retarded time-varying Green's Function solution
 - Average over molecules in 2.5 nm volume to get Macroscopic Maxwell Equations

Maxwell Equations

- Maxwell put Div D into continuity equation
- Added term is displacement current
- Boundary conditions are same as in electro- and magnetostatics

$$\nabla \cdot \bar{J} + \frac{\partial \rho}{\partial t} = \nabla \cdot \left(\bar{J} + \frac{\partial \bar{D}}{\partial t} \right) = 0$$

$$\rightarrow \rightarrow \rightarrow \rightarrow$$

$$\nabla \cdot \bar{D} = \rho$$

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}$$

$$\nabla \cdot \bar{B} = 0$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$$

Auxiliary Mathematical Potentials now Both Scalar and Vector

- $\text{Div } \mathbf{B} = 0$ allows \mathbf{B} to be represented by $\text{curl } \mathbf{A}$
- $\text{Curl } (\mathbf{E} \text{ plus time derivative of } \mathbf{A}) = 0$ says that this quantity can be described by Gradient Φ
- Thus both \mathbf{A} and Φ are required.

$$\nabla \cdot \bar{\mathbf{B}} = 0$$

$$\bar{\mathbf{B}}(\bar{\mathbf{x}}) = \nabla \times \bar{\mathbf{A}}(\bar{\mathbf{x}})$$

$$\nabla \times \left(\bar{\mathbf{E}} + \frac{\partial \bar{\mathbf{A}}}{\partial t} \right) = 0$$

$$\bar{\mathbf{E}} = -\nabla \Phi - \frac{\partial \bar{\mathbf{A}}}{\partial t}$$

Wave Equations for Auxiliary Mathematical Potentials \mathbf{A} and Φ under Lorentz Gauge

$$\nabla^2 \Phi + \frac{\partial}{\partial t} (\bar{A}) = -\rho / \epsilon_0$$

$$\nabla^2 \bar{A} - \frac{1}{c^2} \frac{\partial^2 \bar{A}}{\partial t^2} - \nabla \left(\nabla \cdot \bar{A} + \frac{1}{c^2} \frac{\partial \Phi}{\partial t} \right) = -\mu_0 \bar{J}$$

$$\nabla^2 \Phi - \frac{1}{c^2} \frac{\partial^2 \Phi}{\partial t^2} = -\rho / \epsilon_0$$

$$\nabla^2 \bar{A} - \frac{1}{c^2} \frac{\partial^2 \bar{A}}{\partial t^2} = -\mu_0 \bar{J}$$

- Lorentz Condition (Gauge) when term in this bracket is zero
- Potentials are still not unique

Coulomb Potential and Transverse Current for Auxiliary Mathematical Potentials $\bar{\mathbf{A}}$ and Φ under $\mathbf{Div} \mathbf{A} = 0$

$$\nabla \cdot \bar{\mathbf{A}} = 0$$

- Coulomb Gauge $\mathbf{Div} \mathbf{A} = 0$

$$\nabla^2 \Phi = -\rho / \epsilon_0$$

- Φ is instantaneous near field

$$\nabla^2 \bar{\mathbf{A}} - \frac{1}{c^2} \frac{\partial^2 \bar{\mathbf{A}}}{\partial t^2} = -\mu_0 \bar{\mathbf{J}} + \frac{1}{c^2} \nabla \frac{\partial \Phi}{\partial t}$$

$$\bar{\mathbf{J}} = \bar{\mathbf{J}}_l + \bar{\mathbf{J}}_t$$

- Longitudinal and transverse current

$$\nabla^2 \bar{\mathbf{A}} - \frac{1}{c^2} \frac{\partial^2 \bar{\mathbf{A}}}{\partial t^2} = -\mu_0 \bar{\mathbf{J}}_t$$

- Only transverse current radiates
- Balloon coated with charge and with oscillating radius does not radiate

Green's Function for the Wave Equation Based on Fourier Representation

$$\nabla^2 \Psi - \frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = -4\pi f(\bar{x}, t)$$

$$\Psi(\bar{x}, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \Psi(\bar{x}, \omega) e^{-i\omega t} d\omega$$

$$f(\bar{x}, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}, \omega) e^{-i\omega t} d\omega$$

$$\Psi(\bar{x}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \Psi(\bar{x}, t) e^{i\omega t} dt$$

$$f(\bar{x}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}, t) e^{i\omega t} dt$$

- Typical Wave Equation
- Fourier Representation
- Fourier Spectrum
- Apply to both source distribution and unknown function

Green's Function for the Wave Equation

Radial Behavior

$$(\nabla^2 + k^2)\Psi(\bar{x}, \omega) = -4\pi f(\bar{x}, \omega)$$

$$(\nabla^2 + k^2)G(\bar{x}, \omega) = -4\pi\delta(\bar{x} - \bar{x}')$$

$$\frac{1}{R} \frac{d^2}{dR^2} (RG_k) + k^2 G_k = -4\pi\delta(\bar{R})$$

$$\frac{d^2}{dR^2} (RG_k) + k^2 RG_k = 0$$

$$RG_k(R) = Ae^{ikR} + Be^{-ikR}$$

$$G_k^\pm(R) = \frac{e^{\pm ikR}}{R}$$

- Typical Wave Equation
- Green's Function
- Boundary free case can only depend on R
- Diff Eq. for R variation
- Normalized
- Outward is -kR

Green's Function for the Wave Equation

Time-Retarded (for propagation)

$$(\nabla^2 + k^2)G(\bar{x}, \omega) = -4\pi\delta(\bar{x} - \bar{x}')\delta(t - t')$$

$$\delta(t - t') \rightarrow e^{-i\omega t'}$$

$$G^\pm(\bar{x}, t, \bar{x}', t') = \frac{\delta\left(t' - \left[t \mp \frac{|\bar{x} - \bar{x}'|}{c}\right]\right)}{|\bar{x} - \bar{x}'|}$$

- Put in time delta
 - Fourier Transform
 - $\tau = t - t'$
 - Transform back
 - Use + for source generation
- This time retardation applies to the vector potential under both the Lorentz and Coulomb Gauges and to the scalar potential only under the Lorenz Gauge

Retarded Solutions for B and E

- Apply operators to get E and B
 - Jefimenko expressions
- Work out the retarded derivatives
 - Heaviside-Feynman expressions
 - $E \sim q\{\text{radial}\}$
 - $B \sim q\{v \text{ cross } \mathbf{R}\}$

Derivation of Macroscopic Equations

- Average over space and time
 - Volume 2.5 nm on a side 1000 atoms
 - Time longer than dielectric relaxation time 10^{-14} s
- Procedure
 - Microscopic equations
 - Tapered support or finite support
 - Free and Bound charges
 - Molecular multipole moments
 - Equal to a collection of point multipoles

Macroscopic Maxwell Equations

- Same Equations
- Constitutive Relationships
- Propagation parameter

$$\nabla \cdot \bar{D} = \rho$$

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t}$$

$$\nabla \cdot \bar{B} = 0$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$$

$$\bar{D} = \epsilon \bar{E}$$

$$\bar{B} = \mu \bar{H}$$

$$k = \omega \sqrt{\mu / \epsilon}$$