

EE243 Advanced Electromagnetic Theory**Lec # 8: Magnetostatics**

- Basic Observations of Magnetic Fields and Forces
- Vector Potential
- Magnetic Moment Density (magnetization)
- Force, Torque, Energy
- Macroscopic Equations and Boundary Conditions
- Applications

Reading: Jackson Ch 5

Overview

- Since there are no magnetic charges, one might think Magnetostatics would be easier.
- But no magnetic charges has the opposite effect of greatly complicating the physics and the math.
 - Start with dipole (higher order) effects
 - Now think in terms of fields that circle rather than diverge (lots of cross products)
 - The potential which for electrostatics was a scalar with a clear physical interpretation becomes a multi-component vector without an intuitive interpretation

Basic Observations

- Statics means charge does not change with time => $\text{div } \mathbf{J} = 0$
- Measure torque
- Biot and Savart Law (1820)
- Generalize for charge in motion
- Force on current element

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0$$

$$\rightarrow_{\text{STATICS}} \nabla \cdot \mathbf{J} = 0$$

$$\mathbf{N} = \mu \times \mathbf{B}$$

$$d\vec{B} = \frac{\mu_0}{4\pi} I \frac{d\vec{l} \times \vec{x}}{|\vec{x}|^3}$$

$$\vec{B} = \frac{\mu_0}{4\pi} q \frac{\vec{v} \times \vec{x}}{|\vec{x}|^3}$$

$$d\vec{F} = I_1 (d\vec{l}_1 \times \vec{B})$$

Basic Observations

- B field from a long wire at distance R
- Force on a closed current loop due to another closed current loop
- Force on parallel wires repel if opposite direction
- Force on a current distribution
- Torque on a current distribution

$$|\vec{B}| = \frac{\mu_0}{4\pi} I \int_{-\infty}^{\infty} \frac{dl}{(R^2 + l^2)^{3/2}} = \frac{\mu_0}{2\pi} \frac{I}{R}$$

$$\vec{F}_{12} = -\frac{\mu_0}{4\pi} I_1 I_2 \oint \oint \frac{d\vec{l}_1 \cdot d\vec{l}_2 \vec{x}_{12}}{|\vec{x}_{12}|^3}$$

$$\frac{dF}{dl} = \frac{\mu_0}{2\pi} \frac{I_1 I_2}{d}$$

$$\vec{F} = \int \vec{J}(\vec{x}) \times \vec{B}(\vec{x}) d^3x$$

$$\vec{N} = \int \vec{x} \times (\vec{J} \times \vec{B}) d^3x$$

Differential Form of Magnetostatics

- Start from current element and integrate
- Get curl integral
- Implies $\text{div } \mathbf{B} = 0$
- Go back to integral and take curl
- Curl $\mathbf{B} = \mu_0 \mathbf{J}$
- Ampere's Law for loop of B field = μ_0 current

$$\vec{B}(\vec{x}) = \frac{\mu_0}{4\pi} \int \vec{J}(\vec{x}') \times \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^3} d^3x'$$

$$\vec{B}(\vec{x}) = \frac{\mu_0}{4\pi} \nabla \times \int \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \frac{\mu_0}{4\pi} \nabla \times \nabla \times \int \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \int \vec{J} \cdot \hat{n} da = \mu_0 I$$

Auxiliary Mathematical Potential now Vector

- Div $\mathbf{B} = 0$ allows \mathbf{B} to be represented by curl \mathbf{A}
- Previous slide shows an \mathbf{A}
- But arbitrary grad ψ can be added
- Choose Div $\mathbf{A} = 0$ Coulomb Gauge

$$\nabla \cdot \vec{B} = 0$$

$$\vec{B}(\vec{x}) = \nabla \times \vec{A}(\vec{x})$$

$$\vec{A}(\vec{x}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x' + \nabla \Psi(\vec{x})$$

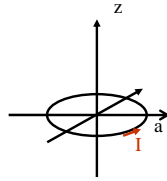
$$\nabla \cdot \vec{A} = 0 \rightarrow \nabla^2 \Psi(\vec{x}) = 0$$

$$\rightarrow \Psi(\vec{x}) = C$$

$$\nabla^2 \vec{A} = -\mu_0 \vec{J}$$

Vector Potential Example: Circular Loop

- Current I on circular loop of radius a
- Integrate J to get A
 - => A_ϕ only
- Apply curl to get B
- Look at far fields
 - Dipole in nature
 - Magnitude $m = \pi a^2 I =$ area times I



$$\vec{A}(\vec{x}) = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{x}}{|\vec{x}|^3}$$

$$\vec{B}(\vec{x}) = \frac{\mu_0}{4\pi} \left[\frac{3\hat{n}(\hat{n} \cdot \vec{m}) - \vec{m}}{|\vec{x}|^3} \right]$$

Copyright 2006 Regents of University of California

Force, Torque and Potential Energy

- Force
 - Example: Magnetic Mirror $\vec{F} = \nabla(\vec{m} \cdot \vec{B})$
- Potential $U = -\vec{m} \cdot \vec{B}$
- Torque
 - Example: Motors $\vec{N} = -\vec{m} \times \vec{B}(0)$

Copyright 2006 Regents of University of California

8

Material Model

- Circular currents due to charge particles orbiting nucleus
 - Sum over particles
 - Related to angular momentum
 - Also quantum effects
 - Local magnetization

$$\vec{m} = \frac{1}{2} \sum_i q_i (\vec{x}_i \times \vec{v}_i)$$

$$\vec{L}_i = M_i (\vec{x}_i \times \vec{v}_i)$$

$$\vec{m} = \frac{e}{2M} \vec{L}$$

$$\vec{M}(\vec{x}) = \sum_i N_i \langle \vec{m}_i \rangle$$

Copyright 2006 Regents of University of California

9

Microscopic to Macroscopic

- Average of $\text{div } \vec{B} = 0$ scales $\nabla \cdot \vec{B} = 0$
- Vector potential A has magnetization contributions (dipole like) $\vec{A}(\vec{x}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{x}') + \nabla' \times \vec{M}(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$
- Magnetization contributes an effective current density $\vec{J} = \nabla \times \vec{M}$
 $\nabla \times \vec{B} = \mu_0 [\vec{J} + \nabla \times \vec{M}]$
- Introduce new macroscopic magnetic field $\vec{H} = \frac{1}{\mu_0} \vec{B} - \vec{M}$
- Curl $\vec{H} = \vec{J}$ $\nabla \times \vec{H} = \vec{J}$

Copyright 2006 Regents of University of California

10

Boundary Conditions

$$(\vec{B}_2 - \vec{B}_1) \cdot \hat{n}_{21} = 0$$

$$\hat{n}_{21} \times (\vec{H}_2 - \vec{H}_1) = \vec{K}_{\text{SURFACE}}$$

$$\hat{n}_{21} = \text{from } 1 \text{ to } 2$$

- Derived from $\text{div } \vec{B} = 0$ and $\text{curl } \vec{H} = \vec{J}$
 - Very large values of permeability occur in which case
 - H is nearly normal to the surface where the B field emerges from the material
 - H is very small inside the material
- Example: electromagnet with N turn coil; Integral of H in loop = NI; and H times gap ~ NI

Copyright 2006 Regents of University of California

11

Boundary Value Problem Techniques

- Choice of convenient Auxiliary function to represent B
- Vector Potential A and $\vec{B} = \text{Curl } \vec{A}$
 - Integrate the vector current with 1/distance to get vector A; take curl to get B; divide μ to get B; apply boundary conditions
- Magnetic Scalar Potential Φ and $\vec{H} = -\text{Grad } \Phi$
 - Special case of no local currents to integrate
 - $\text{Curl } \vec{H} = 0 \Rightarrow \vec{H} = -\text{Grad } \Phi$
 - Choose potential; Grad to get H, multiply μ to get B; apply boundary conditions

Copyright 2006 Regents of University of California

12

Example BVP in Magnetostatics

- Uniformly magnetized sphere
 - m = volume times M
- Magnetized sphere in and external field
 - Analogous to electrostatics
 - Internal Magnetization
- Magnetic Shielding
 - High μ material attracts nearly all field lines
- Hole in a conducting plane in uniform magnetic field
 - Dipole in direction of field

Faraday's Law of Induction

- Flux through surface $F = \int_S \vec{B} \cdot \hat{n} da$
- Electro Motive $Emf = \oint_C \vec{E}' \cdot d\vec{l}$
Force that opposes is proportion to rate of change $Emf = -\frac{dF}{dt}$
- Include velocity of moving loop $\oint_C [\vec{E}' - (\vec{v} \times \vec{B})] \cdot d\vec{l} = -\oint_S \frac{\partial \vec{B}}{\partial t} \cdot \hat{n} da$
- Differential form $\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$

Energy in Magnetic Fields

- Power = VI
- Small loop $\frac{dW}{dt} = -IEmf = I \frac{dF}{Dt}$ $W = \frac{1}{2} I \Delta F$
- A and J in space $\delta W = I \delta F$ $W = \frac{1}{2} \int \vec{H} \cdot \vec{B} d^3x$
- B and H in space (most physical) $\delta W = J \Delta \sigma \int \hat{n} \cdot \delta \vec{B} d^3x$ $W = \frac{1}{2} \int \vec{J} \cdot \vec{A} d^3x$
- Force is gradient of W at constant current $\delta W = \int \delta \vec{A} \cdot \vec{J} d^3x$ $W = \frac{1}{2} \int \vec{M} \cdot \vec{B} d^3x$
 $\vec{F} = (\nabla W)_j$

Self- and Mutual Inductances

- Define using analogy to capacitance $W = \frac{1}{2} \sum_{i=1}^N L_i I_i^2 + \sum_{i=1}^N \sum_{j>i}^N M_{ij} I_i I_j$
- Sort formula for energy into current loops $W = \frac{\mu_0}{8\pi} \int d^3x \int d^3x' \frac{\vec{J}(\vec{x}) \cdot \vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|}$
- When loop small compared to change in vector potential over it the flux linkage can be used $L_{ij} = \frac{\mu_0}{4\pi a^2} \int_{C_i} d^3x_i \int_{C_j} d^3x_j \frac{\vec{J}(\vec{x}_i) \cdot \vec{J}(\vec{x}_j)}{|\vec{x}_i - \vec{x}_j|}$
 $M_{ij} = \frac{\mu_0}{4\pi I_i I_j} \int_{C_i} d^3x_i \int_{C_j} d^3x_j \frac{\vec{J}(\vec{x}_i) \cdot \vec{J}(\vec{x}_j)}{|\vec{x}_i - \vec{x}_j|}$
 $M_{ij} = \frac{\mu_0}{I_i} F_{ij}$