

Similarity transforms App E.13:  $\mathbf{x} = P\bar{\mathbf{x}}$ .

Inverted pendulum

$$ml^2\ddot{\theta} - mgl \sin \theta = T \quad (1)$$

Linearized transfer function:

$$\frac{\Theta(s)}{U(s)} = \frac{1}{s^2 - g/l} = \frac{1}{(s + \Omega)(s - \Omega)} = \quad (2)$$

In state space form with compensator  $D(s) = \frac{U(s)}{U_1(s)} = \frac{s - \hat{\Omega}}{s}$  :

$$\dot{u}(t) = \dot{u}_1(t) - \hat{\Omega}u_1(t) \quad (3)$$

By integration of eqn( 3) we get

$$u(t) = u_1(t) - \hat{\Omega} \int_0^t u_1(\tau) d\tau \quad (4)$$

Introduce state variable  $x_3 = \hat{\Omega} \int u_1$ , then  $\dot{x}_3 = \hat{\Omega}u_1(t)$  and  $u = u_1 - x_3$ . Thus state space representation with compensation is:

$$\dot{\mathbf{x}} = A\mathbf{x} + Bu = \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 \\ \Omega^2 & 0 & -1 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \\ \hat{\Omega} \end{bmatrix} u_1(t) \quad (5)$$

Eigenvalues  $|\lambda I - A| = 0$ :

$$\begin{vmatrix} \lambda & -1 & 0 \\ -\Omega^2 & \lambda & 1 \\ 0 & 0 & \lambda \end{vmatrix} = \lambda(\lambda^2 - \Omega^2) = 0 \quad (6)$$

Thus eigenvalues are  $\lambda_0 = 0$ ,  $\lambda_1 = \Omega$ , and  $\lambda_2 = -\Omega$ . and eigenvectors  $(\lambda I - A)\mathbf{e}_i = \mathbf{0}$ :

$$\mathbf{e}_0 = \begin{bmatrix} 1 \\ 0 \\ \Omega^2 \end{bmatrix} \quad \mathbf{e}_1 = \begin{bmatrix} 1 \\ \Omega \\ 0 \end{bmatrix} \quad \mathbf{e}_2 = \begin{bmatrix} 1 \\ -\Omega \\ 0 \end{bmatrix} \quad (7)$$

To diagonalize  $A$ , let  $\bar{A} = P^{-1}AP$ , where

$$P = [\mathbf{e}_1 \ \mathbf{e}_2 \ \mathbf{e}_0] = \begin{bmatrix} 1 & 1 & 1 \\ \Omega & -\Omega & 0 \\ 0 & 0 & \Omega^2 \end{bmatrix} \quad \text{and} \quad P^{-1} = \frac{1}{2\Omega^2} \begin{bmatrix} \Omega^2 & \Omega & -1 \\ \Omega^2 & -\Omega & -1 \\ 0 & 0 & 2 \end{bmatrix} \quad (8)$$

Thus

$$\bar{A} = P^{-1}AP = \frac{1}{2\Omega^2} \begin{bmatrix} \Omega^2 & \Omega & -1 \\ \Omega^2 & -\Omega & -1 \\ 0 & 0 & 2 \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 \\ \Omega^2 & 0 & -1 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 1 & 1 \\ \Omega & -\Omega & 0 \\ 0 & 0 & \Omega^2 \end{bmatrix} = \begin{bmatrix} \Omega & 0 & 0 \\ 0 & -\Omega & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad (9)$$

Output

$$\mathbf{y} = C\mathbf{x} = CP\bar{\mathbf{x}} = [\mathbf{1} \ \mathbf{0} \ \mathbf{0}] \begin{bmatrix} 1 & 1 & 1 \\ \Omega & -\Omega & 0 \\ 0 & 0 & \Omega^2 \end{bmatrix} \bar{\mathbf{x}} = [\mathbf{1} \ \mathbf{1} \ \mathbf{1}]\bar{\mathbf{x}} \quad (10)$$

$$(11)$$

A variation of the modal canonical form with  $\mathbf{y} = C\mathbf{x}$  is shown here:

